

## Mechanical cleaning of graphene

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Contamination of graphene due to residues from nanofabrication often introduces background doping and reduces electron mobility. For samples of high electronic quality, post-lithography cleaning treatments are therefore needed. We report that mechanical cleaning based on contact mode atomic force microscopy removes residues and significantly improves the electronic properties. A mechanically cleaned dual-gated bilayer graphene transistor with hexagonal boron nitride dielectrics exhibited a mobility of  $\sim 36\,000\text{ cm}^2/\text{Vs}$  at low temperature. © 2012 American Institute of Physics. [doi:10.1063/1.3685504]

High electronic quality is demanded for many graphene experiments<sup>1,2</sup> but is not easily realized. Graphene samples for electronic measurements are typically made with lithographic methods. Lithography makes a myriad of devices possible but always leaves resist residues behind. Making contacts to graphene with shadow mask evaporation solves this contamination issue, but this method has many drawbacks concerning the flexibility of the fabrication process. Hence cleaning after lithography is a crucial step towards obtaining high electronic quality samples. There are different methods at hand: chemical cleaning,<sup>3</sup> thermal cleaning (annealing in an oven),<sup>4,5</sup> and current-induced cleaning.<sup>6,7</sup> Each of these can be very useful but has its own limitations.

In this paper, we present an alternative cleaning method: mechanical cleaning. Scanning a contact mode atomic force microscope (CM AFM) tip over a graphene surface removes residues, removes doping, and improves the electronic mobility without damaging the graphene.

We demonstrate the effectiveness of this method for 4 bilayer graphene on hexagonal boron nitride (hBN) samples. hBN flakes are deposited by mechanical exfoliation on silicon wafers coated with a silicon oxide ( $\text{SiO}_2$ ) layer of thickness  $t_{\text{SiO}_2} = 285\text{ nm}$ . On top of the hBN, we transfer a bilayer graphene flake using a dry transfer method following the protocol of Ref. 8 (at a temperature of  $100^\circ\text{C}$  to remove any water absorbed on the surface of the graphene and hBN flakes). Samples are subsequently annealed in an oven at  $400^\circ\text{C}$  (Ar 2400 sccm,  $\text{H}_2$  700 sccm) to remove residues induced by the transfer process. Cr/Au electrodes are fabricated using electron-beam lithography. We annealed the samples again (same flow rate as the first annealing step) to remove fabrication residues. Trying to clean the graphene, we performed multiple annealing steps at temperatures from  $300^\circ\text{C}$  to a maximum of  $440^\circ\text{C}$ .

After the final annealing step, the samples were often still contaminated. The tapping mode AFM (TM AFM) image of sample A (Fig. 1) shows lots of deposited material

outside the marked window. The roughness in this area is  $\sim 1\text{ nm}$ . Before lithography, all samples were almost atomically flat with a roughness of at most  $0.2\text{ nm}$  (limited by the resolution of the AFM).

We characterize the electronic quality of the samples via the residual doping and field effect mobility. Fig. 2 (lower trace) shows the conductivity of sample B as a function of the backgate voltage at room temperature (RT) in vacuum. From this and similar traces for the other devices, we extracted the charge neutrality point ( $V_{np}$ ) and mobility ( $\mu$ ) of the samples. Depending on the device, we were able to do 2, 3, or 4 terminal measurements. In the two and three terminal measurements, contact resistances make the mobility appear lower. To calculate the conductivity from the measured conductance, we need the aspect ratio of the devices. As some devices were not rectangular, the aspect ratio was hard to determine and we used an underestimated value in

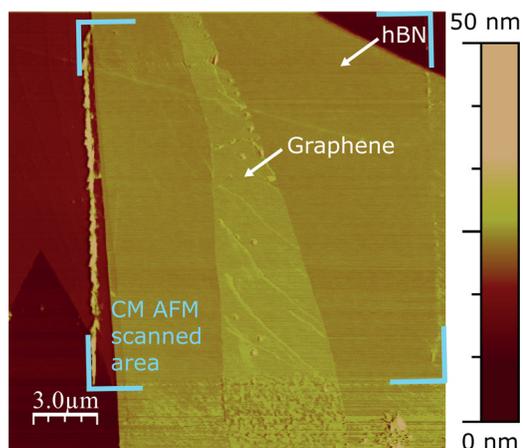


FIG. 1. (Color online) Tapping mode image of sample A after annealing at  $440^\circ\text{C}$  and contact mode scanning (both with a Veeco Nanoscope IIIa AFM). Only the part within the marked window was scanned with the CM AFM. We chose to show this device because it was much more contaminated than other devices before scanning, so that the effect of the CM AFM scan is easily visible. Wrinkles and some tears on the upper right side of the graphene are induced by the tip but were not observed in other devices. On the left and right of the bounding box, walls of deposited residue are visible. The contacts of the device are not visible in this image.

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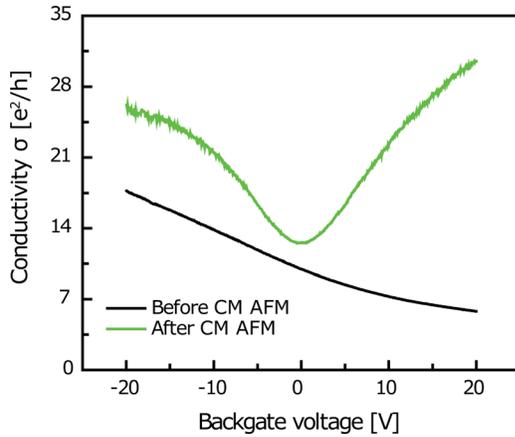


FIG. 2. (Color online) (a) Backgate traces of sample B at room temperature in vacuum ( $I_{bias} = 100$  nA). The lower curve is before CM AFM imaging and the upper curve after. Using the geometric capacitance, we convert the backgate axis into carrier density. Then we extract the field effect mobility by fitting a straight line to the steepest part of the backgate trace:  $\mu = \left( \frac{\epsilon_{SiO_2}}{\epsilon_0 \epsilon_{r,SiO_2}} + \frac{t_{hBN}}{\epsilon_0 \epsilon_{r,hBN}} \right) \frac{d\sigma}{dV}$ , where  $\epsilon_{r,SiO_2} = 3.9$  and  $\epsilon_{r,hBN} = 3.0$ , as calculated from Fig. 3. We extracted the thickness  $t_{hBN}$  from TM AFM images of the devices.

our calculations. Those two factors make the measured mobility a lower bound of the actual mobility. The results for all samples are summarized in Table I. The electronic measurements indicate that most of the devices were highly doped and had a relatively low mobility. Most likely residues on top of the graphene as seen in the AFM images induce doping and provide scattering centers that degrade the electronic quality.<sup>9</sup>

To remove the residues, we scanned the samples in contact mode AFM with a constant force (Veeco OTR8-35 tip with a stiffness of 0.15 N/m). Hereby, the tip is held in contact with the sample surface. We engaged the tip with the lowest force possible. When the tip made contact, we confirmed a reasonable set-point force with the help of a force distance measurement, discussed further below. Then we started scanning the sample with a rate of 0.5–1 Hz. For most samples, we scanned the same area several times but without further visible effect.

Tapping mode images taken after scanning in CM AFM show that we cleaned the graphene (Fig. 1 inside marked window). The roughness is at most 0.2 nm, similar to the values measured before processing the devices. Further evidence that we removed residues from the graphene are the banks of deposits that are visible in Fig. 1, exactly at the boundaries of the area that was scanned in contact mode.

After CM and TM AFM imaging, we again recorded backgate traces at room temperature in vacuum (upper curve

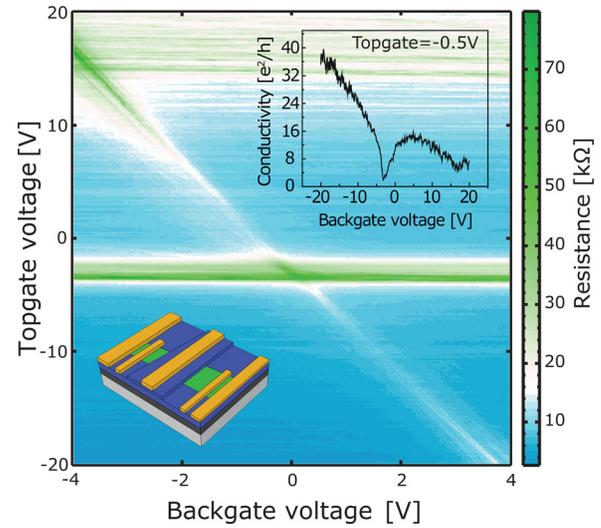


FIG. 3. (Color online) Measurements on a double gated bilayer graphene transistor fabricated out of sample A. The 4-probe resistance at  $T = 50$  mK is plotted as a function of backgate and topgate voltage ( $V_{bg}$  and  $V_{tg}$ , respectively). From the slope of the diagonal line, we calculated the relative dielectric constant of the hBN to be 3.0 assuming a parallel plate capacitor model and  $\epsilon_{r,SiO_2} = 3.9$ . The thickness of the bottom hBN flake was 14 nm and the top hBN flake 50 nm, values extracted from AFM images. Lowerleft inset: schematic of the device. Blue colored regions are hBN, green is bilayer graphene, and yellow are the contacts and gate. Upper right inset: resistance as a function of  $V_{bg}$  at  $V_{tg} = -0.5$  V. The dip at  $V_{bg} \sim 17$  V is caused by the uncovered graphene part.

in Fig. 2). Not only the mobility increased twofold but also doping was reduced. For other samples, we observed similar behavior (see Table I).

An attractive feature of mechanical cleaning is that it can be naturally followed by further sample processing. We fabricated sample A into a double gated bilayer device.<sup>10,11</sup> With the same dry transfer method as mentioned before, we stamped an hBN flake on sample A that will act as a topgate dielectric. We defined a topgate electrode across the flake and two voltage probes by e-beam lithography (lower left inset Fig. 3). Resistance as a function of the topgate and backgate voltages is plotted in Fig. 3. The resistance peaks that run horizontally are caused by the regions next to the topgated region. The diagonal ridge shows the typical increase of resistance due to opening of a gap with increasing perpendicular electric field.<sup>10</sup> From the upper right inset, we extract a (hole) mobility  $\mu$  of  $\sim 36000$  cm<sup>2</sup>/Vs at carrier density  $n \sim 5 \times 10^{10}$  cm<sup>-2</sup> (corrected for the change in slope due to the neutrality point around 17 V). This value is among the highest found in the literature for bilayer graphene devices, including suspended devices.<sup>8,12,13</sup>

We now turn to the mechanism by which CM AFM removes residues from the sample surface. Presumably, the

TABLE I. Results summary for four different samples (measurements in vacuum).

Sample	Anneal T (°C)	No. of passes	Scan force (nN)	Measurement	$V_{np}$ before (V)	$V_{np}$ after (V)	$\mu$ before (cm <sup>2</sup> /Vs)	$\mu$ after (cm <sup>2</sup> /Vs)
A	440	6	2.3	2-prb, RT	4	-7	$3.4 \times 10^3$	$8.9 \times 10^3$
A	440	6	2.3	4-prb, 50mK	n.a.	-3	n.a.	$3.6 \times 10^4$
B	360	1	-2.9	3-prb, RT	>20	0	$1.7 \times 10^3$	$2.8 \times 10^3$
C	360	2	-4.6	3-prb, RT	>20	1	$2.6 \times 10^2$	$9.2 \times 10^2$
D	360	5	-22	4-prb, RT	17	-1	$2.7 \times 10^3$	$6.7 \times 10^3$

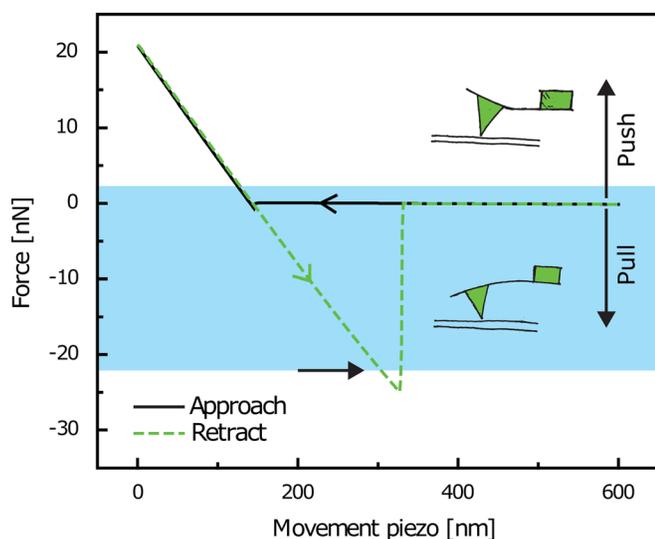


FIG. 4. (Color online) Force-distance curve of sample D, measured by holding the tip of the AFM in a fixed lateral position and approaching and retracting the tip in the vertical direction. While making these vertical movements, the deflection of the tip is recorded. Assuming that when the tip is in contact with the surface, the tip deflects the same distance as the piezo moves, we can calibrate the deflection scale. With the spring constant of the tip, we convert that deflection to a force. The horizontal axis has an arbitrary offset. The blue region indicates the range of forces that we used for cleaning the samples. Sample D was scanned at a force of  $-22$  nN as indicated by the arrow. The illustrations picture the pulling and pushing regime.

tip is plowing through a layer of physisorbed contaminants and thereby “brooms” the graphene clean, which means the interaction of the tip with the surface is larger than the interaction of the contaminants with the surface. We believe that removing residues is the main explanation for the improvement of the electronic quality of the graphene. In principle, the CM AFM might also flatten the graphene, reducing ripples, and thereby enhancing mobility. However, we observe no difference in the flatness of graphene before fabrication (presumably equal to the state after lithography) and after mechanical cleaning. Flattening of the graphene should thus play little or no role. The hBN substrate does appear to play a role in improving electronic quality. Jalilian *et al.*<sup>14</sup> deployed the mechanical cleaning on single layer graphene samples on a  $\text{SiO}_2$  substrate. They also observed an improvement in surface morphology, but electronic quality did not improve. We observed the same behavior in a single layer graphene on  $\text{SiO}_2$  sample. Further research needs to be done to explain the role of the substrate.

To gain more insight in the interplay of the surface and the CM AFM tip, we took force-distance curves (Fig. 4) in the area we scanned in contact mode. From these curves, we can extract the force we were exerting on the sample during scanning, which ranged from  $-22$  nN to  $+2.3$  nN depending

on the device. A positive force denotes the tip was pushing and a negative force denotes the tip was pulling on the surface. In pulling configuration, the tip is held in contact by the Van der Waal’s interaction and adhesive forces due to water. The broad range of scanning forces that gave good results illustrates the robustness of the mechanical cleaning method.

In summary, scanning bilayer graphene on hBN in CM AFM removes contaminants from the surface, reduces residual doping, and significantly improves electronic mobility. A double gated bilayer graphene transistor, which was mechanically cleaned in the fabrication process, showed mobilities up to  $36\,000\text{ cm}^2/\text{Vs}$  at 50 mK and opening of a bandgap. This illustrates the effectiveness and versatility of CM AFM for obtaining high-quality graphene devices. Possibly, AFM and scanning tunneling microscope (STM) setups in vacuum could benefit even more from mechanical cleaning as it can be applied *in situ*, avoiding subsequent contamination by molecules absorbed from the air.<sup>4,15</sup>

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- <sup>1</sup>A. H. C. Neto, F. Guinea, N. M. R. Peres, K. S. Novoselov, and A. K. Geim, *Rev. Mod. Phys.* **81**, 109 (2009).
- <sup>2</sup>M. S. Fuhrer, C. N. Lau, and A. H. MacDonald, *MRS Bull.* **35**, 289 (2010).
- <sup>3</sup>Z. Cheng, Q. Zhou, C. Wang, Q. Li, C. Wang, and Y. Fang, *Nano Lett.* **11**, 767 (2011).
- <sup>4</sup>M. Ishigami, J. H. Chen, W. G. Cullen, M. S. Fuhrer, and E. D. Williams, *Nano Lett.* **7**, 1643 (2007).
- <sup>5</sup>Y. Dan, Y. Lu, N. J. Kybert, Z. Luo, and A. T. C. Johnson, *Nano Lett.* **9**, 1472 (2009).
- <sup>6</sup>J. Moser, A. Barreiro, and A. Bachtold, *Appl. Phys. Lett.* **91**, 163513 (2007).
- <sup>7</sup>K. Bolotin, K. Sikes, Z. Jiang, M. Klima, G. Fudenberg, J. Hone, P. Kim, and H. Stormer, *Solid State Commun.* **146**, 351 (2008).
- <sup>8</sup>C. Dean, A. Young, I. Meric, C. Lee, L. Wang, S. Sorgenfrei, K. Watanabe, T. Taniguchi, P. Kim, K. Shepard, and J. Hone, *Nat. Nanotechnol.* **5**, 722 (2010).
- <sup>9</sup>J. Chen, C. Jang, S. Adam, M. S. Fuhrer, E. D. Williams, and M. Ishigami, *Nat. Phys.* **4**, 377 (2008).
- <sup>10</sup>J. B. Oostinga, H. B. Heersche, X. Liu, A. F. Morpurgo, and L. M. K. Vandersypen, *Nature Mater.* **7**, 151 (2008).
- <sup>11</sup>Y. Zhang, T. Tang, C. Girit, Z. Hao, M. C. Martin, A. Zettl, M. F. Crommie, Y. R. Shen, and F. Wang, *Nature* **459**, 820 (2009).
- <sup>12</sup>B. E. Feldman, J. Martin, and A. Yacoby, *Nat. Phys.* **5**, 889 (2009).
- <sup>13</sup>A. S. Mayorov, D. C. Elias, M. Mucha-Kruczynski, R. V. Gorbachev, T. Tudorovskiy, A. Zhukov, S. V. Morozov, M. I. Katsnelson, V. I. Fal’ko, A. K. Geim *et al.*, *Science* **333**, 860 (2011).
- <sup>14</sup>R. Jalilian, L. A. Jauregui, G. Lopez, J. Tian, C. Roecker, M. M. Yazdanpanah, R. W. Cohn, I. Jovanovic, and Y. P. Chen, *Nanotechnology* **22**, 295705 (2011).
- <sup>15</sup>A. Castellanos-Gomez, N. Agrait, and G. Rubio-Bollinger, *Nanotechnology* **21**, 145702 (2010).