

The cost of destroying entanglement and resource

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Outline for section 1

- 1 The problem of destroying correlation
- 2 The problem of destroying entanglement
- 3 Transforming a state to locally recovered states
- 4 Generalization to other resource theories
- 5 Techniques

The set up



ρ_{AB}



Groisman, Popescu, Winter [Phys. Rev. A., 2005]

The set up


 ρ_{AB}
 $\sigma_{A'}$


$$\tau_J = \frac{1}{N} \sum_{j=1}^N |j\rangle\langle j|_J$$

Groisman, Popescu, Winter [Phys. Rev. A., 2005]

The set up


 U_j
 $\sigma_{A'}$
 J
 ρ_{AB}


$$\tau_J = \frac{1}{N} \sum_{j=1}^N |j\rangle\langle j|_J, \quad U_{AA'J} = \sum_j |j\rangle\langle j|_J \otimes U_j$$

Groisman, Popescu, Winter [Phys. Rev. A., 2005]

The set up



$$\rho'_{AA'B}$$



$$\rho'_{AA'B} = \text{Tr}_J(U_{AA'J}\rho_{AB} \otimes \sigma_{A'} \otimes \tau_J U_{AA'J}^\dagger)$$

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The set up



$$\rho'_{AA'B} \approx \rho_A \otimes \rho_B \otimes \sigma_{A'}$$



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- In the asymptotic and i.i.d. setting, with $\rho_{AB}^{\otimes n}$.
- Showed that $\frac{1}{n} \log |J| \rightarrow I(A : B)_\rho$.
- Achievability does not use A' register.
- $I(A : B)_\rho = S(\rho_A) + S(\rho_B) - S(\rho_{AB})$.
- $I(A : B)_\rho = D(\rho_{AB} \| \rho_A \otimes \rho_B) = \inf_{\sigma_A, \sigma_B} D(\rho_{AB} \| \sigma_A \otimes \sigma_B)$.
- Measure of correlation.

Outline for section 2

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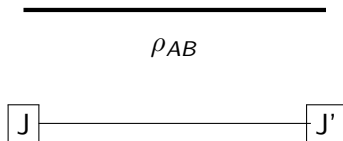
Our set up



ρ_{AB}

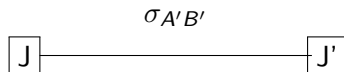
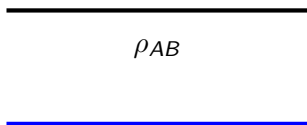


The set up



$$\tau_{JJ'} = \frac{1}{N} \sum_{j=1}^N |j\rangle\langle j|_J \otimes |j\rangle\langle j|_{J'}$$

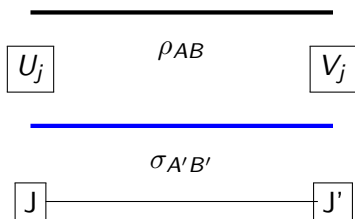
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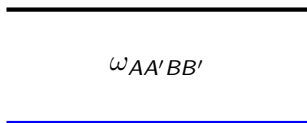
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$$\omega_{AA'BB'} = \text{Tr}_{JJ'}((U_{AA'J} \otimes V_{BB'J'}) \rho_{AB} \otimes \sigma_{A'B'} \otimes T_{JJ'} (U_{AA'J}^\dagger \otimes V_{BB'J'}^\dagger))$$

The set up



$$\omega_{AA'BB'} \approx \theta_{AB} \otimes \sigma_{A'B'}$$



$$\omega_{AA'BB'} = \text{Tr}_{JJ'}((U_{AA'J} \otimes V_{BB'J'}) \rho_{AB} \otimes \sigma_{A'B'} \otimes \tau_{JJ'} (U_{AA'J}^\dagger \otimes V_{BB'J'}^\dagger))$$

θ_{AB} is separable

Some definitions

- Definitions:

$$E(\rho_{AB}) = \inf_{\sigma_{AB} \in \text{SEP}(A:B)} D(\rho_{AB} \parallel \sigma_{AB})$$

$$E_{\varepsilon}^{\text{one-shot}}(\rho_{AB}) = \inf_{\sigma_{AB} \in \text{SEP}(A:B)} D_{\max}^{\varepsilon}(\rho_{AB} \parallel \sigma_{AB})$$

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- $E(\rho_{AB})$: relative entropy of entanglement. (Vedral, Plenio, Rippin, Knight [Phys. Rev. Lett. 1997]).
 - An entanglement measure.
 - Monotonic under non-entangling operations.
 - Natural interpretation under quantum hypothesis testing (Brandao, Plenio [Comm. Math. Phys., 2010]).

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- $D_{\max}(\rho \parallel \sigma) = \inf \{k : \rho \preceq 2^k \sigma\}$ (Datta [IEEE TIT, 2009]).
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- Robustness of entanglement (Vidal, Tarrach [Phys. Rev. A., 1999]).

Our results

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Theorem (AA/MH/RJ; MB/CM)

There exists a protocol with error $\varepsilon + \delta$ requires discarding $E_{\varepsilon}^{\text{one-shot}}(\rho_{AB}) + 2 \log \frac{1}{\delta}$ bits. Furthermore any protocol with error ε must discard $E_{\varepsilon}^{\text{one-shot}}(\rho_{AB})$ bits.

- Error measured by the purified distance (Gilchrist, Langford, Nielsen [Phys. Rev. A., 2005]).

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- In the asymptotic and i.i.d. setting, this becomes $\lim_{n \rightarrow \infty} \frac{1}{n} E(\rho_{AB}^{\otimes n})$.
- Recall: $E(\rho_{AB}) = \inf_{\sigma_{AB} \in \text{SEP}(A:B)} D(\rho_{AB} \| \sigma_{AB})$.
- Thus, disentanglement cost is an entanglement measure.

Outline for section 3

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Locally recovered states

- Fix a tripartite quantum state ρ_{ABC} .
- The set of all locally recovered states are $(I_B \otimes \mathcal{R}_{C \rightarrow AC})(\rho_{BC})$.
- Relative entropy of recovery:

$$D(A : B|C)_\rho = \inf_{\mathcal{R}_{C \rightarrow AC}} D(\rho_{ABC} \| (I_B \otimes \mathcal{R}_{C \rightarrow AC})(\rho_{BC})).$$

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- Regularized version is $D^\infty(A : B|C)_\rho$.
 - A natural interpretation as hypothesis testing between a quantum state and the set of locally recovered states. (Cooney, Hirche, Morgan, Olson, Seshadreesan, Watrous, Wilde [Phys. Rev. A., 2016]).
 - A lower bound on the quantum conditional mutual information. (Brandao, Harrow, Oppenheim, Winter [Phys. Rev. Lett., 2015]).

Transformation via local mixtures of unitaries

- Local mixture of unitaries is of the following form

$$\frac{1}{J} \sum_{i=1}^J U_A^i \otimes V_B^i \otimes W_C^i(\cdot) U_A^{i\dagger} \otimes V_B^{i\dagger} \otimes W_C^{i\dagger}.$$

- How much noise ($\log J$) is required to convert ρ_{ABC} to a locally recovered state, allowing catalysts that are locally recovered states?

Characterization of the noise

- Achievability and a converse:

Theorem (MB/CM)

There exists a protocol that requires a rate of $D^\infty(A : B|C)_\rho$ bits in the asymptotic and i.i.d. setting. Furthermore any protocol for which the unitary V_B^i are permutations requires a rate of $D^\infty(A : B|C)_\rho$ bits.

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- The achievability protocol above has the property that V_B^i are permutation operations.
- Open question to find near optimal characterization under arbitrary local mixture of unitaries.

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- Free states given by \mathcal{F} . We assume:
 - Convex set.
 - Closed under tensor product.
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$$E(\rho) = \inf_{\sigma \in \mathcal{F}} D(\rho || \sigma).$$

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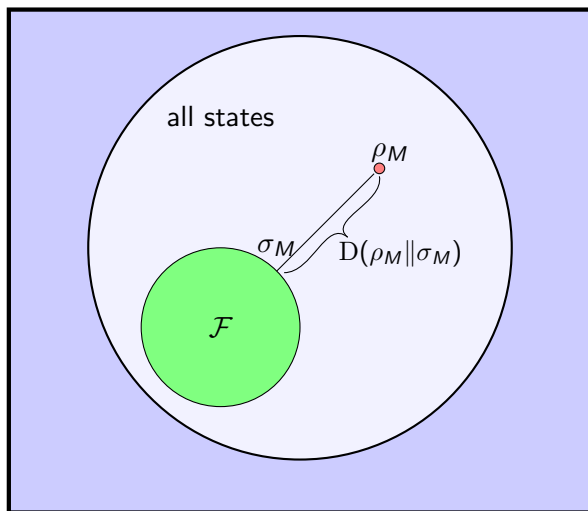
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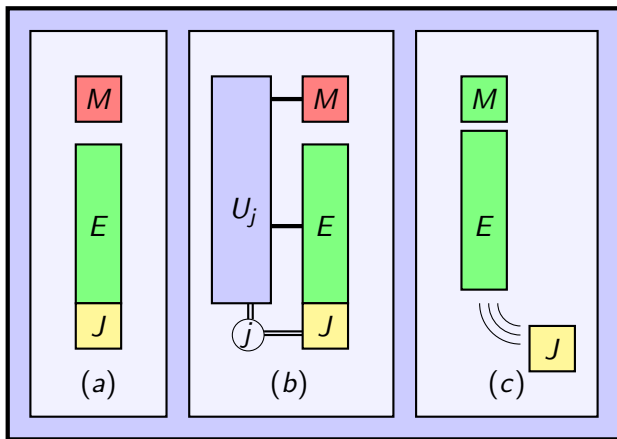
- Max-relative entropy of resource:

$$E_{\varepsilon}^{\text{one-shot}}(\rho) = \inf_{\sigma \in \mathcal{F}} D_{\max}^{\varepsilon}(\rho \| \sigma)$$

Relative entropy of resource: geometric view



Our task



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The role of relative entropy of resource

- Brandao and Gour [Phys. Rev. Lett., 2015] show that in the asymptotic and i.i.d. setting, the rate of transformation of ρ to σ is

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- Brandao and Gour [Phys. Rev. Lett., 2015] show that in the asymptotic and i.i.d. setting, the rate of transformation of ρ to σ is

$$\frac{E(\rho)}{E(\sigma)}.$$

- Suggests that $E(\rho)$ is a natural measure for quantifying the amount of resource in a state.

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- Coherence.
 - Braumgatz, Cramer, Plenio [Phys. Rev. Lett. 2014]; Winter, Yang [Phys. Rev. Lett., 2015]; Streltsov, Adesso, Plenio [Rev. Mod. Phys., 2017].
- Asymmetry.
 - Wakakuwa [Phys. Rev. A., 2017], where relative entropy of frameness is non-zero. Does not apply to the formulation in Gour, Marvian, Spekkens [Phys. Rev. A., 2009].

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- Quantum thermodynamics.
 - Brandao et. al. [Phys. Rev. Lett., 2013]; Brandao et. al. [PNAS, 2015]; Horodecki, Oppenheim [Nat. Comm., 2013]; Faist et. al. [Nat. Comm., 2015]; Gour et. al. [Phys. Rep. 2015]; Narasimhachar, Gour [Nat. Comm., 2015].

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- The converse proofs follow from basic one-shot entropic inequalities adapted to respective settings.

Convex-split Lemma

- A., Devabathini, Jain [Phys. Rev. Lett. 2017, arXiv 2014].
- Let Ψ_{RB}, σ_B be quantum states, $k = D_{\max}(\Psi_{RB} \| \Psi_R \otimes \sigma_B)$.

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- Consider the following quantum state

$$\tau_{RB_1 B_2 \dots B_N} = \frac{1}{N} \sum_{j=1}^N \Psi_{RB_j} \otimes \sigma_{B_1} \otimes \sigma_{B_2} \dots \otimes \sigma_{B_{j-1}} \otimes \sigma_{B_{j+1}} \dots \otimes \sigma_{B_N}$$

Convex-split Lemma

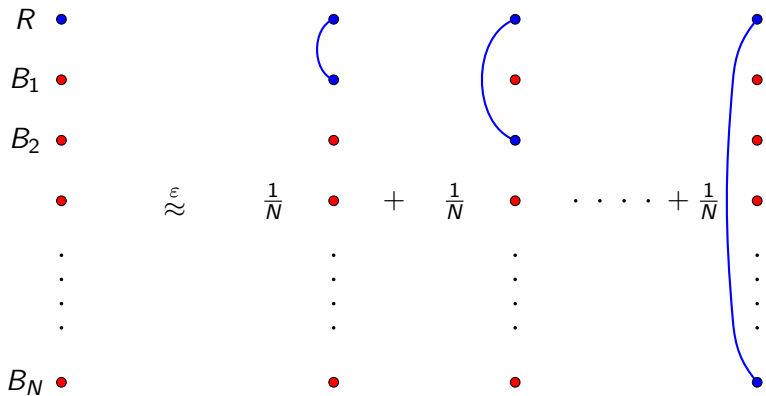
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- Then,

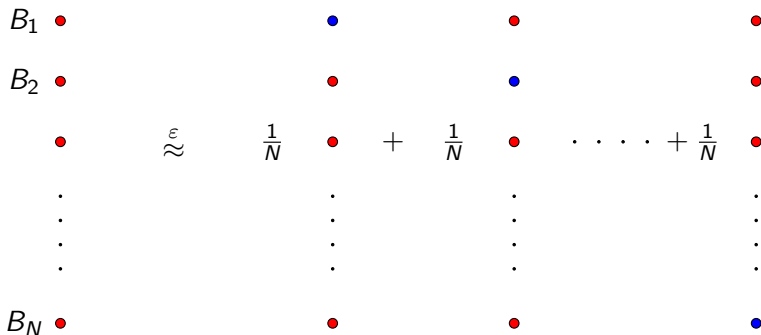
$$D(\tau_{RB_1 B_2 \dots B_N} \| \Psi_R \otimes \sigma_{B_1} \otimes \sigma_{B_2} \dots \otimes \sigma_{B_N}) \leq \frac{2^k}{N}.$$

Convex-split Lemma: In pictures



If $\log N \geq D_{\max}(\Psi_{RB} \| \Psi_R \otimes \sigma_B) + \log \frac{1}{\epsilon}$.

Convex-split Lemma: for our application



$$\text{If } \log N \geq D_{\max}(\Psi_B \parallel \sigma_B) + \log \frac{1}{\epsilon}.$$

Protocol

ρ_M ●

σ_{M_1} ●

σ_{M_2} ●

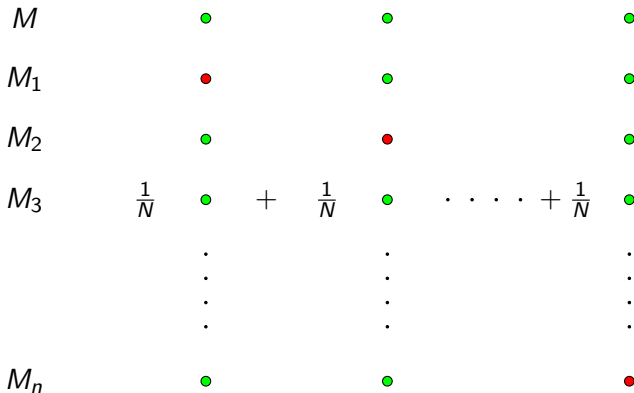
⋮

σ_{M_n} ●

σ is the state that minimizes

$$E_\varepsilon^{\text{one-shot}}(\rho) = \inf_{\sigma \in \mathcal{F}} D_{\max}^\varepsilon(\rho \| \sigma).$$

Protocol



Conclusion

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- Obtains near optimal bounds in settings we consider.

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- The achievability protocol uses simple random swap operation.
- Obtains near optimal bounds in settings we consider.
- Open: reduce the amount of catalyst required (compare randomness extraction).
- Open: one shot analogue for resource transformation (Brandao, Gour [Phys. Rev. Lett., 2015]).

Thank you for your attention!